

# Entanglement and Trainability

## Variational Quantum Algorithms · Module 4 — Barren Plateaus

### Node 4.3

#### Position in Knowledge Graph

**System:** Variational Quantum Algorithms **Structural Domain:** Barren Plateaus **Node:** 4.3

The original barren plateau theorem (4.2) is phrased in terms of unitary 2-designs and Haar averaging. This is mathematically clean but operationally opaque — “the circuit forms a 2-design” doesn’t tell you what the *states* look like or what it would take to avoid the plateau.

This node gives a **second, physically transparent view** of the plateau: barren plateaus are what you get when the ansatz produces **highly entangled** states. More specifically, when an  $n$ -qubit ansatz reaches states with **volume-law** entanglement across any bipartition, the reduced density matrix of any subsystem is close to maximally mixed, and any local observable has a vanishing expectation and vanishing gradient.

This is the **Marrero-Kieferová-Wiebe (2021)** framing: barren plateaus and volume-law entanglement are two sides of the same coin.

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#### Entanglement Scaling Laws

To set up the argument, recall the basic scaling laws for entanglement of pure  $n$ -qubit states.

Let  $|\psi\rangle$  be an  $n$ -qubit pure state, and let  $A$  be a subsystem of  $n_A \leq n/2$  qubits. Compute the reduced density matrix  $\rho_A = \text{Tr}_{\bar{A}}|\psi\rangle\langle\psi|$ . The entanglement entropy of  $\rho_A$  is

$$S(\rho_A) = -\text{Tr}(\rho_A \log \rho_A).$$

**Three regimes, set by how  $S(\rho_A)$  depends on  $n_A$ :**

- **Area law:**  $S(\rho_A) = O(\text{boundary area of } A)$ . For a 1D system, this is  $S(\rho_A) = O(1)$  — constant in subsystem size. Ground states of gapped local Hamiltonians satisfy area laws.
- **Logarithmic law:**  $S(\rho_A) = O(\log n_A)$ . Ground states of gapless local Hamiltonians (critical points).
- **Volume law:**  $S(\rho_A) = O(n_A)$  — linear in subsystem size. Haar-random states, thermal states at high temperature, states generated by random deep circuits.

The volume-law regime is the “highly entangled” regime. It saturates the maximum possible entanglement:  $S(\rho_A) \leq n_A \cdot \log 2$ , with equality when  $\rho_A = I/2^{n_A}$  (maximally mixed).

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#### Volume Law Implies Plateau

**Theorem (Marrero, Kieferová, Wiebe, informal):** If the ansatz  $U(\theta)$  produces volume-law entangled states for typical  $\theta$ , then:

1. The reduced density matrix  $\rho_A(\theta)$  is exponentially close to the maximally mixed state  $I/2^{n_A}$ .

2. Any local observable  $O_A$  supported on  $A$  has expectation value  $\langle O_A \rangle = \text{Tr}(O_A)/2^{n_A}$ , which for traceless  $O_A$  is zero.
3. The variance of  $\langle O_A \rangle$  across random  $\theta$  is exponentially small in  $n_A$ .
4. The variance of the gradient of  $\langle O_A \rangle$  across random  $\theta$  is exponentially small in  $n - n_A$  (the complement size).

This is the barren plateau, rediscovered from a purely state-based argument. No Haar integration needed — just entanglement scaling.

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## Why This Matters: The Tradeoff

This framing makes the barren plateau problem look like a **tradeoff between expressivity and trainability**:

- **High expressivity** = ability to reach many different quantum states = typically implies *generating highly entangled states* = typically volume-law = typically barren plateau.
- **Low expressivity** = restricted set of states = typically lower entanglement = typically trainable.

The more states your ansatz can prepare, the more likely it is that random initialization lands on a highly-entangled state, and the more likely the gradient is unusable.

**This is the central tradeoff of ansatz design.** You can't have “reaches everything” and “easy to train” simultaneously — they are in direct tension.

Module 8 (Expressivity vs Trainability) is dedicated to this tradeoff in full. The point of this node is to introduce the entanglement angle on it, which is both physically clearer and useful for designing mitigations.

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## Entanglement-Aware Ansatz Design

The entanglement view of plateaus directly suggests mitigation strategies that focus on *controlling* entanglement:

- 1. Shallow circuits.** Shallow circuits cannot generate volume-law entanglement. For depth  $d$ , the entanglement of a typical reduced density matrix is  $O(d)$  (each gate can add  $O(1)$  entanglement). So a circuit with  $d \ll n$  has  $O(d)$ -area-law entanglement, not volume-law, and doesn't trigger the entanglement-based plateau.
- 2. Sparse entanglement graphs.** Limit the entanglement-generating gates (CNOTs) to a specific graph. If the graph is sparse (say, a line or a tree), entanglement between distant qubits is slow to build up, and the ansatz stays in a low-entanglement regime.
- 3. Locality-preserving ansätze.** Design gates so that each one only acts on  $O(1)$  qubits. Then in constant depth, the entanglement is area law. HEA with nearest-neighbor coupling satisfies this in the shallow regime.
- 4. Entanglement regularization.** Add a term to the cost function that penalizes excess entanglement (e.g.,  $\mu S(\rho_A)$ ). This pushes the optimizer into the low-entanglement regime, where gradients are usable.
- 5. Matrix product state ansätze.** Use ansätze (like MPS, TTN, PEPS) that are **guaranteed by construction** to have bounded entanglement. These ansätze cannot represent volume-law states at all, so they are trivially plateau-free in the entanglement sense.

These mitigations are all justified by the entanglement view; they would be harder to motivate from the Haar-2-design view alone.

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## The Dual Failure Mode: Not Enough Entanglement

There is a flip side to this story worth flagging, even though Module 4 is mostly about too much entanglement.

If the ansatz produces states with *too little* entanglement, it can't represent the ground state of many interesting Hamiltonians. For example, the ground state of a 1D critical model has logarithmic entanglement, but a product-state ansatz has zero. A product-state ansatz is trivially trainable (no plateau) but represents a much smaller set of states than is needed.

**The sweet spot is: enough entanglement to represent the ground state, but not so much that the plateau engages.**

For quantum chemistry, ground states typically have area-law or weak-entanglement structure, so low-entanglement ansätze can work. For condensed matter and high-energy physics, ground states may have logarithmic or volume-law structure, and the tradeoff is harder to navigate.

This is one of the principal *structural* arguments for why VQAs are more tractable for chemistry than for condensed matter.

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## Entanglement and Shot-Noise Sensitivity

An additional consequence of the entanglement view: **high-entanglement states make the cost estimator noisier.**

For a maximally entangled state, any local Pauli has  $\langle P \rangle = 0$  and variance  $1 - 0 = 1$  per shot. For a low-entanglement product-like state, many local Paulis have  $\langle P \rangle$  near  $\pm 1$ , giving variance near 0.

So the entanglement-induced plateau is not only a gradient problem but a *cost estimation* problem. When the ansatz is in the high-entanglement regime, the cost estimator has maximum variance, *and* the gradient is exponentially small. Both effects compound — you need exponentially more shots to detect an exponentially smaller signal, a double-exponential cost.

This is why the entanglement view is operationally sharper: it predicts not just the gradient scaling but also the shot-cost scaling, and the two combine to make the plateau regime completely uneconomical.

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## Connecting To The 2-Design View

The two views of barren plateaus — “Haar 2-design on parameters” and “volume-law entanglement” — are **equivalent** in the sense that they describe the same physical regime.

Specifically: - A Haar-random unitary  $U$  applied to  $|0\rangle$  produces a Haar-random pure state  $|\psi\rangle$ . - Haar-random pure states have volume-law entanglement with overwhelming probability (Page's theorem). - So any ansatz whose parameter distribution approximates a 2-design produces volume-law entangled states with overwhelming probability. - And vice versa: any ansatz that produces volume-law entangled states for typical parameters satisfies the 2-design assumption well enough for the barren plateau theorem to apply.

The 2-design view is mathematically cleaner for proving the theorem; the entanglement view is physically cleaner for motivating mitigations. Both are correct; they are the same thing told two ways.

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## Structural Role

This node is the **physical interpretation** of the theorem from 4.2. It explains what barren plateaus *look like* in terms of the quantum states being prepared, and sets up the depth-scaling and mitigation discussions in 4.4 and 4.6.

It is also the node that forwards cleanest to Module 8 (Expressivity vs Trainability), which develops the entanglement-expressivity-trainability tradeoff in quantitative form.

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## Relations to Other Concepts

- **Page’s theorem** — Haar-random states have near-maximal entanglement.
  - **Area laws** — entanglement of ground states of gapped local Hamiltonians.
  - **Matrix product states** — tensor network ansätze with bounded entanglement by construction.
  - **Entropy of entanglement** — the von Neumann entropy of the reduced density matrix.
  - **Cerezo et al. (2021)** — the local cost function refinement, which can be viewed as an entanglement-aware fix.
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## Worked Example — Entanglement Of A Volume-Law State

Consider  $n = 10$  qubits, bipartition  $A = \{0, 1, 2, 3, 4\}$  and  $\bar{A} = \{5, 6, 7, 8, 9\}$ .

For a Haar-random state, Page’s theorem gives

$$\mathbb{E}[S(\rho_A)] \approx n_A \log 2 - \frac{2^{n_A-1}}{2^{n-n_A}} = 5 \log 2 - \frac{16}{32} = 5 \log 2 - 0.5 \approx 2.97.$$

The maximum possible entanglement is  $n_A \log 2 \approx 3.47$ , so a Haar-random state is within 0.5 of maximum — essentially saturated.

The reduced density matrix is

$$\rho_A \approx \frac{I_A}{32} + O(2^{-5}).$$

So for any traceless observable  $O_A$  (say  $O_A = Z_0$ ):

$$\langle O_A \rangle = \text{Tr}(\rho_A O_A) \approx 0 + O(2^{-5}) \approx 0.03.$$

The expectation is essentially zero, with fluctuation  $\sim 2^{-5}$  — which matches the barren plateau variance scaling.

At  $n = 20$ , the same calculation gives  $\langle O_A \rangle \sim 2^{-10} \approx 10^{-3}$  — in the unmeasurable regime.

This is the entanglement-based barren plateau: **the state is so scrambled that any local measurement gives a fixed answer (the Haar average), with vanishing variation across parameter changes.**

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## Visualization

*Pending. A bar chart: entanglement entropy  $S(\rho_A)$  on the y-axis, vs ansatz type on the x-axis (product, shallow HEA, deep HEA, Haar-random). A horizontal dashed line marks the volume-law threshold. Below the threshold: gradients trainable. Above: barren plateau. Caption: “Entanglement is the operational axis of the barren plateau.”*

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## Exercises

1. For a single Bell state  $|00\rangle + |11\rangle$  on 2 qubits, compute the reduced density matrix on qubit 0 and verify it is maximally mixed. What is  $\langle Z_0 \rangle$ ?
  2. For a GHZ state on 3 qubits,  $(|000\rangle + |111\rangle)/\sqrt{2}$ , compute the reduced density matrix on the first 2 qubits. Is it maximally mixed? What does this say about the entanglement of GHZ compared to Haar-random?
  3. (VQA) Design an ansatz that provably produces states with at most logarithmic entanglement. What kinds of ground states could such an ansatz reach? What does this imply about its usefulness for chemistry vs condensed matter?
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## Research Extensions

- **Page-curve corrections to plateau bounds** — using the full Page distribution to sharpen the barren plateau variance calculation.
  - **Entanglement-regularized training** — adding entropy penalties to the cost function to keep the optimizer in trainable regions.
  - **MPS-based VQA** — replacing quantum circuits with MPS ansätze (or hybrid quantum-MPS) to guarantee bounded entanglement.
  - **Entanglement monitoring during training** — using intermediate entanglement measurements to detect and exit plateau regimes.
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## Cross-links to Other Courses

- **Quantum Information Theory** — entanglement entropy, Page’s theorem, maximally mixed states.
  - **Condensed Matter Physics** — area laws, topological entanglement, ground-state entanglement.
  - **Tensor Networks** — MPS, PEPS, MERA — classical ansätze with bounded entanglement.
  - **Module 8 (Expressivity vs Trainability)** — the full tradeoff in quantitative form.
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*Variational Quantum Algorithms · Module 4 — Barren Plateaus · Node 4.3*

**Concepts:** barren-plateaus, entanglement, von-neumann-entropy, density-matrix

**Prerequisites:** 4.2

**Enables:** 4.4, 4.6

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